

Discrimination of muons for mass composition studies of inclined air showers detected with IceTop

The IceCube Collaboration

(a complete list of authors can be found at the end of the proceedings)

E-mail: aswathi.balagopalv@icecube.wisc.edu

IceTop, the surface array of IceCube, measures air showers from cosmic rays within the energy range of 1 PeV to a few EeV and a zenith angle range of up to $\approx 36^\circ$. This detector array can also measure air showers arriving at larger zenith angles at energies above 20 PeV. Air showers from lighter primaries arriving at the array will produce fewer muons when compared to heavier cosmic-ray primaries. A discrimination of these muons from the electromagnetic component in the shower can therefore allow a measurement of the primary mass. A study to discriminate muons using Monte-Carlo air showers of energies 20-100 PeV and within the zenith angular range of 45° - 60° will be presented. The discrimination is done using charge and time-based cuts which allows us to select muon-like signals in each shower. The methodology of this analysis, which aims at categorizing the measured air showers as light or heavy on an event-by-event basis, will be discussed.

Corresponding authors: Aswathi Balagopal V.^{1*}

¹ *Wisconsin IceCube Particle Astrophysics Center, University of Wisconsin, Madison, WI 53703, USA*

* Presenter

*37th International Cosmic Ray Conference (ICRC 2021)
July 12th – 23rd, 2021
Online – Berlin, Germany*

1. Introduction

The determination of the primary mass composition of cosmic rays with varying energies provides added information about the origin of the cosmic rays that contribute to the observed cosmic-ray spectrum and their propagation within the Universe. A common method to determine the mass of the cosmic-ray primary is with the use of the muon content within the air shower. Heavy primaries, like iron, will generate a larger number of muons in the air shower than light primaries, like proton, of the same energy.

In this analysis, we explore a method that can be used to determine the light or heavy mass group of the primary cosmic ray on an event-by-event basis. The study is performed for inclined air showers detected by the IceTop array. Such a study is important for the test of existing hadronic interaction models, and the muon content for different mass groups predicted by such models. The composition study for inclined air showers measured by IceTop will also provide an independent cross-check to other composition analyses of IceTop using quasi-vertical air showers. IceTop is the surface component of IceCube, the neutrino telescope at the South Pole. IceTop consists of 162 ice-Cherenkov tanks that measure signals deposited by charged particles arising from cosmic-ray air showers [1]. These 162 tanks are grouped into 81 stations, with each station composed of two tanks. Figure 1 shows the geometry and layout of the IceTop array. Each IceTop tank contains two digital optical modules (DOMs) that collect the deposited Cherenkov light. Since Cherenkov light is emitted by both muons and electromagnetic particles (with non-negligible hadronic contribution), a combination of signals from these will be observed in the tank signal.

There exist several analyses using IceTop and IceCube with the goal of discriminating the electromagnetic particles from the muons. For analyses focused on the determination of the muon density, it is crucial to get a near-perfect number of the observed muons [2]. However, in the cases where the discrimination is used to determine the mass-composition, a contamination of electromagnetic components within the muon parameter can be tolerated [3].

Other analyses using IceTop and IceCube for mass-composition studies operate only for showers with zenith angles up to 36° [3–6]. Air showers with zenith angles (θ) larger than this are generally not included within the standard IceTop analyses. This is primarily due to the unavailability of energy calibration (for the energy proxy S_{125} which is the signal expected to be deposited in a tank at a distance of 125 m) using simulations for these air showers. However, we can determine the energy of these showers with larger θ using the method of constant intensities, where we compare the inclined showers to the vertical showers (with the expectation of universal flux rates) to determine their energy [7]. Therefore, we can use such inclined air showers also for mass-composition studies; which is done in this analysis.

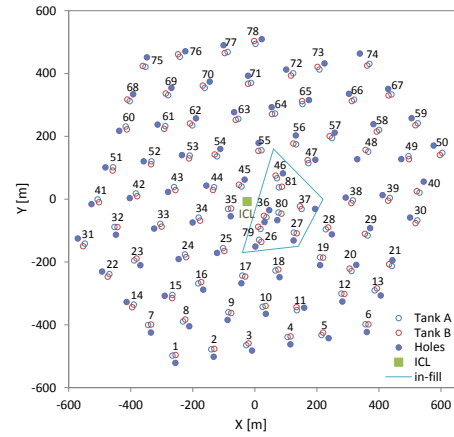


Figure 1: The layout of the IceTop array from [1].

2. Analysis method

We use air-shower simulations that were generated using CORSIKA [8], with Sibyll 2.1 [9] as the hadronic-interaction model. We consider simulations with zenith angles above 45° and energies above $10^{7.3}$ GeV. This is the energy threshold of inclined air showers measured with IceTop [7]. Since the available simulation set (with full detector response) ends at an energy of 10^8 GeV, this is the maximum energy that we consider in this analysis. We divide the showers into three zenith bins: 45° - 50° , 50° - 55° , and 55° - 60° . These showers are also divided into different bins of their energy proxy $\log_{10}(S_{125}/\text{VEM})$, where VEM is the expected amount of charge deposited in a tank by a vertical-equivalent muon. We choose only the bins of $\log_{10}(S_{125}/\text{VEM})$ where the energy is above $10^{7.3}$ GeV. Also, some energy-proxy bins are dropped due to insufficient statistics for the simulations. The energy-proxy bins that we use are: $\log_{10}(S_{125}/\text{VEM}) = 0.7$ - 1.0 , 1.0 - 1.3 , 1.3 - 1.6 for 45° - 50° showers, and $\log_{10}(S_{125}/\text{VEM}) = 0.4$ - 0.7 , 0.7 - 1.0 , 1.0 - 1.3 for showers with zenith angles 50° - 55° and 55° - 60° [7].

We apply cuts to the simulated events to ensure good reconstruction quality for the events used in the analysis. These event-level quality cuts are as follows: 1. We require 5 IceTop stations to be triggered and passed through the filter. 2. The reconstruction of the air shower is required to have succeeded. 3. The number of layers of tanks around the shower core should be greater than one. This ensures that at least two outer-layers of tanks are there surrounding the shower core. 4. The tank with the maximum signal within the shower should not be located at the edge of the array. Even if the previous cut is applied, we sometimes are left with events where the maximum signal is in the edge, due to noisy hits. Application of this cut ensures better quality for the events chosen. Although we reject $\approx 80\%$ events with these cuts, the remaining events have good quality, which is needed to conduct this study.

Apart from these event-level quality cuts, we apply some analysis-level cuts to select the muon-like signals in the air showers. The first cut is a charge-based cut. We choose those hits in the tanks which have a charge between 1 and 4 VEM. This was chosen since muons from inclined air showers were mainly seen to fall within this charge range for IceTop. This cut already removes a majority of the electromagnetic contamination within the selected hits; especially at larger distances from the shower core. In addition to this, we apply a time-based cut to the remaining hits in the event. Only those events with $\Delta(\delta t) > 0$ are chosen in the final level. Here, δt is the curvature of the shower front. $\Delta(\delta t) = \delta t_{\text{observed}} - \delta t_{\text{predicted}}$ compares the observed curvature to the theoretical prediction of the shower curvature (which is optimised for quasi-vertical showers). The condition $\Delta(\delta t) > 0$ chooses the hits that arrived before their predicted times. Combined with the charge-based cut, this allows us to choose the early muons in the shower. The addition of this cut further removes electromagnetic contamination from the selected hits. Figure 2 shows the effect of the analysis-level cuts that selects muon-like hits for proton and iron showers.

Once the final-level hits have been selected, they are binned with respect to their perpendicular distance to the shower axis. The sum of all hits above each distance bin is evaluated to obtain the muon-like parameter.

$$\xi_{\mu\text{-like}} = \sum N_{\mu\text{-like}} > r \quad (1)$$

The lateral distribution of the muon-like parameter, $\xi_{\mu\text{-like}}$ for air showers within the zenith angular range of 55° - 60° for the energy bin $\log_{10}(S_{125}/\text{VEM}) = 0.7$ - 1.0 is shown in Figure 3.

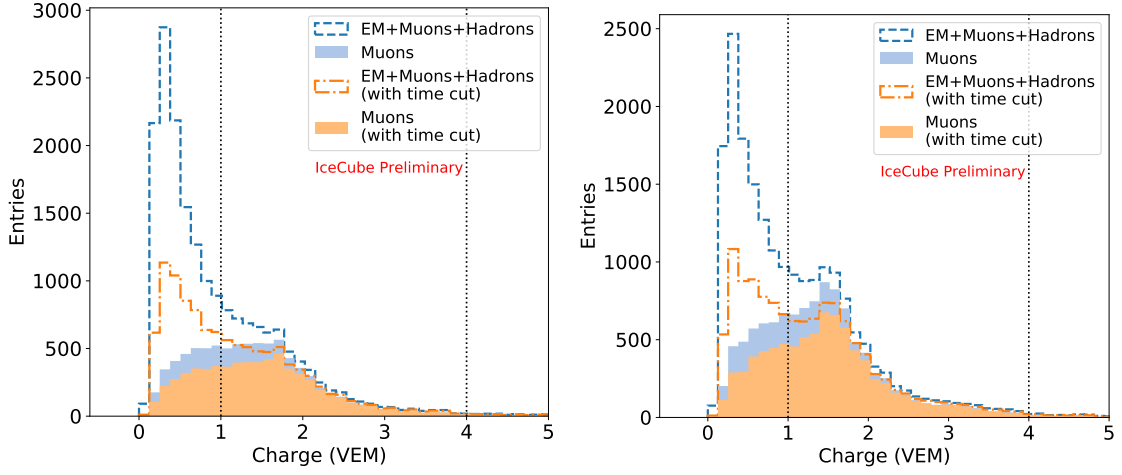


Figure 2: Left: Proton showers, $\theta = 50^\circ - 55^\circ$ and $\log_{10}(S_{125}/\text{VEM}) = 0.7 - 1.0$. The charge within the tanks from all particle types for hits with distance from the shower axis ≥ 260 m are shown. These are the events that pass the event-level quality cuts. The muon-component of these hits are also shown. The charge distribution once the time-based cut is applied is also shown (including its muon content). All hits that fall within the two vertical lines between 1 and 4 VEM will be chosen as the final muon-like hits. Right: The same for iron showers, $\theta = 45^\circ - 50^\circ$ and $\log_{10}(S_{125}/\text{VEM}) = 1.0 - 1.3$. The hits with distance from the shower axis ≥ 330 m are shown.

This is shown for simulated air showers of different primaries. As expected, the mean of $\xi_{\mu\text{-like}}$ exhibits the trend where heavier primaries have more muon-like hits than lighter primaries. We can also see that there is a good amount of separation between the distribution of the light and the heavy primaries. We can therefore use this parameter to determine the lightness/heaviness of each shower. The bottom panel of Figure 3 depicts the fraction of electromagnetic contamination within $\xi_{\mu\text{-like}}$ at each distance bin. The electromagnetic contamination is seen to be small, especially at distances further from the center of the shower.

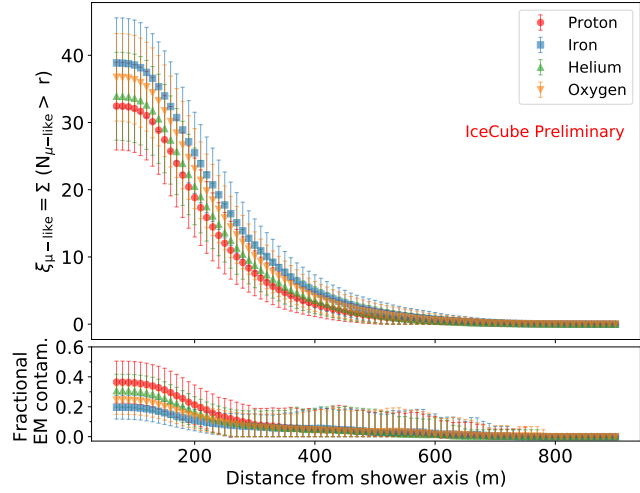


Figure 3: Top panel: lateral distribution of the muon-like parameter, $\xi_{\mu\text{-like}}$, for showers from different primaries with $\theta = 55^\circ - 60^\circ$ and $\log_{10}(S_{125}/\text{VEM}) = 0.7-1.0$. Bottom panel: corresponding fraction of electromagnetic contamination within signals considered in the $\xi_{\mu\text{-like}}$ calculation. Both the top and bottom panels show the mean values (marker) and standard deviations (error bar) of the respective distribution.

The next step in the analysis is the calculation of the reference distance at which we get the maximum separation between the $\xi_{\mu\text{-like}}$ of the heavy and light primaries. For this we look at the proton and iron primaries alone. A Gaussian distribution is fit to the $\xi_{\mu\text{-like}}$ at each distance bin for both the proton and iron showers. A measure of the separation of these two distributions is estimated to get the distance at which the separation is maximum, with least amount of fluctuations.

For this, the figure of merit (FOM) is calculated as

$$\text{FOM} = \frac{|\mu(\text{H}) - \mu(\text{Fe})|}{\sqrt{\sigma_{\text{H}}^2 + \sigma_{\text{Fe}}^2}}, \quad (2)$$

where μ is the mean of the Gaussian distribution and σ is the standard deviation [10].

An example of the FOM calculation is given in Figure 4 (left). The figure shows the distributions of $\xi_{\mu\text{-like}}$ at a distance of 260 m for proton and iron showers with $\theta = 55^\circ - 60^\circ$ and with $\log_{10}(S_{125}/\text{VEM}) = 0.7 - 1.0$. The resulting value of the FOM for this distribution is 1.08 ± 0.08 . Once the FOM is calculated for each distance bin, they can be compared as shown in Figure 4 (right). The figure shows the FOM values at different distances for $\theta = 50^\circ - 55^\circ$ and $\log_{10}(S_{125}/\text{VEM}) = 1.0 - 1.3$. Based on this, the distance of 260 m is chosen as the reference distance for this bin. This procedure is repeated for all zenith bins and all energy bins used in the analysis.

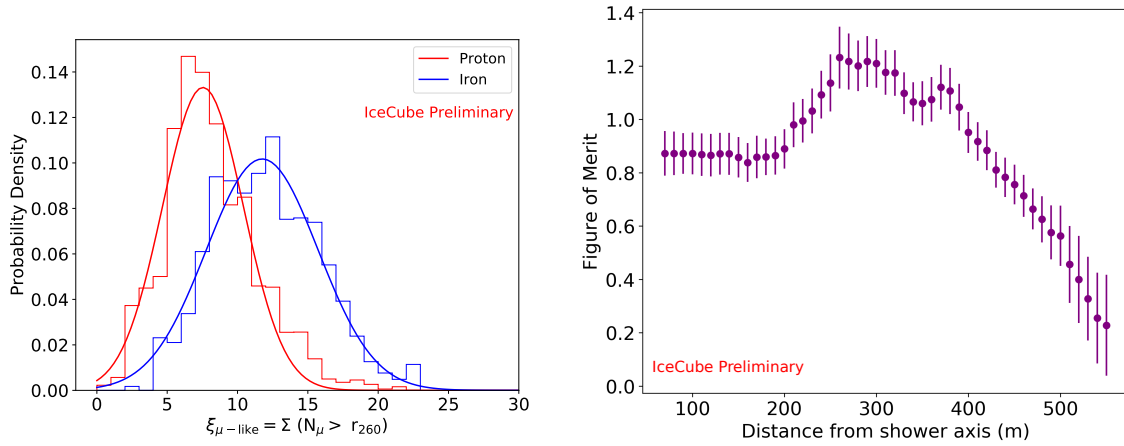


Figure 4: Left: An example of the FOM calculation for showers with $\theta = 55^\circ - 60^\circ$ and $\log_{10}(S_{125}/\text{VEM}) = 0.7 - 1.0$. A Gaussian distribution is fit to both iron and proton distributions of $\xi_{\mu\text{-like}}$ for a given distance bin (here 260 m). Right: The obtained FOM vs distance for showers with $\theta = 50^\circ - 55^\circ$ and $\log_{10}(S_{125}/\text{VEM}) = 1.0 - 1.3$. The reference distance is chosen from this.

Upon the determination of the reference distance, the probability-density distribution for $\xi_{\mu\text{-like}}$ can be drawn for both proton and iron showers. This is done for each zenith and energy bin considered in the analysis. An example of such a distribution for showers with $\theta = 50^\circ - 55^\circ$ and $\log_{10}(S_{125}/\text{VEM}) = 0.7 - 1.0$ is shown in Figure 5. The figure shows the two-dimensional probability-density distribution of $\xi_{\mu\text{-like}}$ at the reference distance of 260 m with respect to $\log_{10}(S_{125}/\text{VEM})$. The solid-blue curves in the figure represent the 1σ , 2σ and 3σ contours for the iron showers while the red-dashed curves represent the same for the proton showers. The right panel in the figure also shows the combined $\xi_{\mu\text{-like}}$ distributions of proton and iron showers within the entire bin. A separation between the iron and proton showers is visible, with some amount of overlap. This signifies that in the regions with less overlap, we can identify the showers more confidently, while in the overlap regions the shower primary becomes more ambiguous.

The nature of the iron and proton probability-density distributions is utilised to evaluate the light or heavy nature of each observed air-shower event. For this, the marginal cumulative

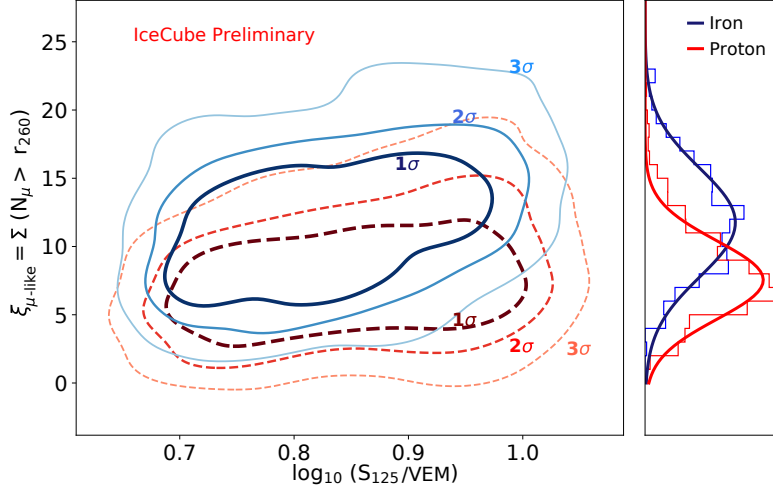


Figure 5: Probability-density distributions for proton (red dashed) and iron (blue solid) showers with $\theta = 50^\circ - 55^\circ$ and $\log_{10}(S_{125}/VEM) = 0.7 - 1.0$ for the reference distance of 260 m. The right panel shows the combined distribution of $\xi_{\mu\text{-like}}$.

distribution (summation along $\xi_{\mu\text{-like}}$ -axis) is determined along each slice of $\log_{10}(S_{125})$. The summation is performed along the positive y-axis for iron and along the negative y-axis for the proton distribution. This results in an iron-marginal cumulative distribution (Fe-CD) value and proton-marginal cumulative distribution (H-CD) value for each point in this space. With this, we can assign a Fe-CD and H-CD value to each event with a given pair of $\log_{10}(S_{125})$ and $\xi_{\mu\text{-like}}$. Any event with a high H-CD value and low Fe-CD value would be identified as light (or proton-like) and vice-versa for heavy (or iron-like) events.

3. Testing the procedure

In order to verify the performance of the analysis procedure, we conduct a test using simulations. The simulated air showers are passed through the analysis to see the fraction of proton/iron events correctly identified as light/heavy.

Figure 6 shows the verification done for iron (left) and proton (right) simulations, on an event-by-event basis. Each simulated event is assigned an H-CD and Fe-CD value, and a histogram of these assigned values is shown in the figure. The colour scale depicts the fraction of events entering each bin in the histogram. Events with low Fe-CD values are considered as light and events with low H-CD values are considered as heavy.

We can sum up the fractions in the bottom-first row of bins on the x-axis (with low H-CD value) to determine the percentage of entries identified as heavy. The same can be done for each set of rows. Similarly, the columns can be added up to see the percentage of entries in each column, the leftmost column (with low Fe-CD) being the events most confidently identified as light.

Table 1 shows the percentage of events that enter the bottom-first row (with low H-CD value) and the bottom-second row, that are identified as heavy for showers with $\theta = 50 - 55^\circ$ and $\log_{10}(S_{125}/VEM) = 0.7-1.0$ (middle θ and energy bins). The table also shows the percentage of

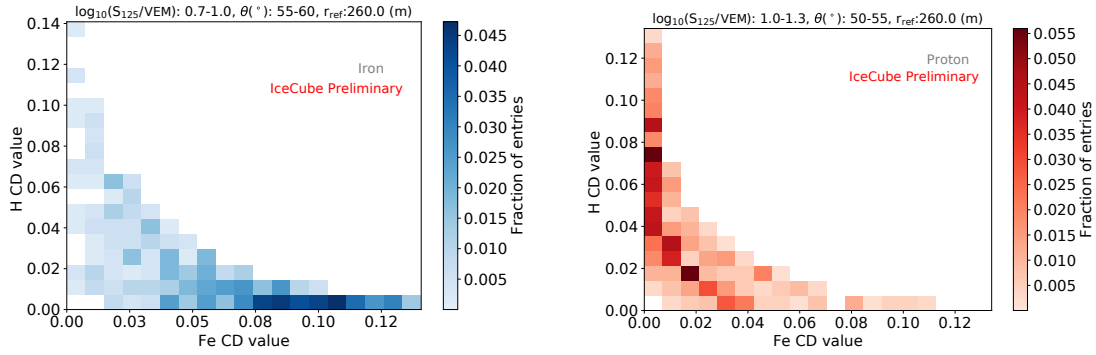


Figure 6: Verification of the mass-classification procedure using simulations. The H-CD value and Fe-CD value is created for each event and compared. Left: Iron showers ($\log_{10}(S_{125}/VEM) = 0.7-1.0$ and $\theta = 55 - 60^\circ$), all events falling in the bins with low H-CD values are correctly identified as heavy. 45.4% events fall in the 1st row (correctly identified) and 2.7% events fall in the 1st column (misidentified as light). Right: Proton showers ($\log_{10}(S_{125}/VEM) = 1.0-1.3$ and $\theta = 50-55^\circ$), with events falling in the bins with low Fe-CD values identified as light. 46.2% events fall in the 1st row (correctly identified) and 12.9% events fall in the 1st column (misidentified as heavy)

Element	Events in 1st column (light)	Events in 2nd column (light)	Events in 1st row (heavy)	Events in 2nd row (heavy)
True H	26.7%	14.6%	6.2%	7.1%
True Fe	3.3%	4.4%	40.2%	17.1%
True He	14.7%	10.6%	16.1%	11.6%
True O	7.1%	6.2%	21.0%	13.6%

Table 1: Testing the classifier for various elements for events that fall in the middle θ and energy proxy bins $\theta = 50 - 55^\circ$ and $\log_{10}(S_{125}/VEM) = 1.0-1.3$.

events in the leftmost column (with low Fe-CD value) and the second column from the left, that are identified as light. The same calculation can be performed for the events shown in Figure 6 also.

It is clear that 40 – 50% of iron events are correctly identified as heavy (with better confidence) and 30 – 40% of proton events are correctly identified as light. A smaller percentage of proton/iron events are incorrectly identified as heavy/light. One caveat to this is that the same events used for drawing the probability-density distributions have been used to test the procedure, due to the lack of sufficient simulations.

The same procedure can be repeated for the available helium and oxygen simulations to test the classifier. This is shown in Figure 7 for helium and oxygen events with $\theta = 45^\circ - 50^\circ$ and $\log_{10}(S_{125}/VEM) = 0.7 - 1.0$. The helium and oxygen events are not seen to clump as closely to the light/heavy composition bins as the proton/iron events. They show mixed nature, as expected. This shows that the classification is indeed composition sensitive since helium and oxygen should fall as intermediaries. The percentage of helium and oxygen events that fall in the light-identified columns and heavy-identified rows for air showers falling in the middle θ and S_{125} bins are shown in Table 1.

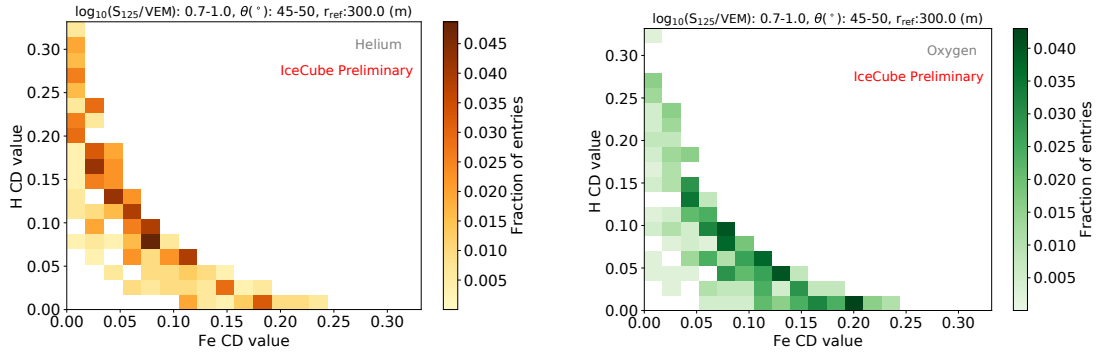


Figure 7: The verification done for helium showers (left) and oxygen showers (right) with, $\log_{10}S_{125}/VEM = 0.7 - 1.0$ and $\theta = 45 - 50^\circ$, using the H-CD and Fe-CD values obtained from the probability-density distributions obtained from proton and iron primaries. 10.0% He events fall in the 1st row (heavy) and 17.2% events fall in the 1st column (light). 20.4% He events fall in the 1st row (heavy) and 7.8% events fall in the 1st column (light).

4. Summary and Outlook

We have explored a new method for determining the composition of inclined air showers on an event-by-event basis. This utilizes the discrimination of muon-like signals in the air shower and using it to draw probability-density distributions for both proton and iron showers. Based on this, we classify events as light or heavy by assigning them marginal cumulative distribution values and marginal inverse cumulative distribution values. A verification of the procedure reveals that the method is effective. Further tests on other hadronic interaction models will be needed to determine the model-dependent performance of the analysis. With that, we can use the analysis method to determine the mass composition of inclined air showers measured with IceTop, and thereby provide an independent check to the composition measurement of quasi-vertical showers from the same detector array.

References

- [1] IceCube Collaboration, R. Abbasi *et al.* *Nucl. Instr. and Meth. in Phys. Res. A* **700** (Feb., 2013) 188–220.
- [2] IceCube Collaboration *PoS ICRC2021* (these proceedings) 342.
- [3] IceCube Collaboration, M. G. Aartsen *et al.* *Phys. Rev. D* **100** no. 8, (2019) 082002.
- [4] IceCube Collaboration *PoS ICRC2021* (these proceedings) 312.
- [5] IceCube Collaboration *PoS ICRC2021* (these proceedings) 323.
- [6] IceCube Collaboration *PoS ICRC2021* (these proceedings) 357.
- [7] A. Balagopal V., [DOI: 10.5445/IR/1000091377](https://doi.org/10.5445/IR/1000091377). PhD thesis, 2019.
- [8] D. Heck, J. Knapp, J. N. Capdevielle, G. Schatz, and T. Thouw. (1998) FZKA-6019.
- [9] E.-J. Ahn *et al.* *Phys. Rev. D* **80** (Nov, 2009) 094003.
- [10] E. M. Holt, F. G. Schröder, and A. Haungs *Eur. Phys. J. C* **79** no. 5, (2019) 371.

Full Author List: IceCube Collaboration

R. Abbasi¹⁷, M. Ackermann⁵⁹, J. Adams¹⁸, J. A. Aguilar¹², M. Ahlers²², M. Ahrens⁵⁰, C. Alispach²⁸, A. A. Alves Jr.³¹, N. M. Amin⁴², R. An¹⁴, K. Andeen⁴⁰, T. Anderson⁵⁶, G. Anton²⁶, C. Argüelles¹⁴, Y. Ashida³⁸, S. Axani¹⁵, X. Bai⁴⁶, A. Balagopal V.³⁸, A. Barbano²⁸, S. W. Barwick³⁰, B. Bastian⁵⁹, V. Basu³⁸, S. Baur¹², R. Bay⁸, J. J. Beatty^{20,21}, K.-H. Becker⁵⁸, J. Becker Tjus¹¹, C. Bellenghi²⁷, S. BenZvi⁴⁸, D. Berley¹⁹, E. Bernardini^{59,60}, D. Z. Besson^{34,61}, G. Binder^{8,9}, D. Bindig⁵⁸, E. Blaufuss¹⁹, S. Blot⁵⁹, M. Boddenberg¹, F. Bontempo³¹, J. Borowka¹, S. Böser³⁹, O. Botner⁵⁷, J. Böttcher¹, E. Bourbeau²², F. Bradascio⁵⁹, J. Braun³⁸, S. Bron²⁸, J. Brostean-Kaiser⁵⁹, S. Browne³², A. Burgman⁵⁷, R. T. Burley², R. S. Busse⁴¹, M. A. Campana⁴⁵, E. G. Carnie-Bronca², C. Chen⁶, D. Chirkin³⁸, K. Choi⁵², B. A. Clark²⁴, K. Clark³³, L. Classen⁴¹, A. Coleman⁴², G. H. Collin¹⁵, J. M. Conrad¹⁵, P. Coppin¹³, P. Correa¹³, D. F. Cowen^{55,56}, R. Cross⁴⁸, C. Dappen¹, P. Dave⁶, C. De Clercq¹³, J. J. DeLaunay⁵⁶, H. Dembinski⁴², K. Deoskar⁵⁰, S. De Ridder²⁹, A. Desai³⁸, P. Desiati³⁸, K. D. de Vries¹³, G. de Wasseige¹³, M. de With¹⁰, T. DeYoung²⁴, S. Dharani¹, A. Diaz¹⁵, J. C. Díaz-Vélez³⁸, M. Dittmer⁴¹, H. Dujmovic³¹, M. Dunkman⁵⁶, M. A. DuVernois³⁸, E. Dvorak⁴⁶, T. Ehrhardt³⁹, P. Eller²⁷, R. Engel^{31,32}, H. Erpenbeck¹, J. Evans¹⁹, P. A. Evenson⁴², K. L. Fan¹⁹, A. R. Fazely⁷, S. Fiedlschuster²⁶, A. T. Fienberg⁵⁶, K. Filimonov⁸, C. Finley⁵⁰, L. Fischer⁵⁹, D. Fox⁵⁵, A. Frankowiak^{11,59}, E. Friedman¹⁹, A. Fritz³⁹, P. Fürst¹, T. K. Gaisser⁴², J. Gallagher³⁷, E. Ganster¹, A. Garcia¹⁴, S. Garrappa⁵⁹, L. Gerhardt⁹, A. Ghadimi⁵⁴, C. Glaser⁵⁷, T. Glauch²⁷, T. Glüsenskamp²⁶, A. Goldschmidt⁹, J. G. Gonzalez⁴², S. Goswami⁵⁴, D. Grant²⁴, T. Grégoire⁵⁶, S. Griswold⁴⁸, M. Gündüz¹¹, C. Günther¹, C. Haack²⁷, A. Hallgren⁵⁷, R. Halliday²⁴, L. Halve¹, F. Halzen³⁸, M. Ha Minh²⁷, K. Hanson³⁸, J. Hardin³⁸, A. A. Harnisch²⁴, A. Haungs³¹, S. Hauser¹, D. Hebecker¹⁰, K. Helbing⁵⁸, F. Henningsen²⁷, E. C. Hettinger²⁴, S. Hickford⁵⁸, J. Hignight²⁵, C. Hill¹⁶, G. C. Hill², K. D. Hoffman¹⁹, R. Hoffmann⁵⁸, T. Hoinka²³, B. Hokanson-Fasig³⁸, K. Hoshina^{38,62}, F. Huang⁵⁶, M. Huber²⁷, T. Huber³¹, K. Hultqvist⁵⁰, M. Hünnefeld²³, R. Hussain³⁸, S. In⁵², N. Iovine¹², A. Ishihara¹⁶, M. Jansson⁵⁰, G. S. Japaridze⁵, M. Jeong⁵², B. J. P. Jones⁴, D. Kang³¹, W. Kang⁵², X. Kang⁴⁵, A. Kappes⁴¹, D. Kappesser³⁹, T. Karg⁵⁹, M. Kar²⁷, A. Karle³⁸, U. Katz²⁶, M. Kauer³⁸, M. Kellermann¹, J. L. Kelley³⁸, A. Kheirandish⁵⁶, K. Kin¹⁶, T. Kintscher⁵⁹, J. Kiryluk⁵¹, S. R. Klein^{8,9}, R. Koirala⁴², H. Kolanoski¹⁰, T. Kontrimas²⁷, L. Köpke³⁹, C. Kopper²⁴, S. Kopper⁵⁴, D. J. Koskinen²², P. Koundal³¹, M. Kovacevich⁴⁵, M. Kowalski^{10,59}, T. Kozynets²², E. Kun¹¹, N. Kurahashi⁴⁵, N. Lad⁵⁹, C. Lagunas Gualda⁵⁹, J. L. Lanfranchi⁵⁶, M. J. Larson¹⁹, F. Lauber⁵⁸, J. P. Lazar^{14,38}, J. W. Lee⁵², K. Leonard³⁸, A. Leszczyńska³², Y. Li⁵⁶, M. Lincetto¹¹, Q. R. Liu³⁸, M. Liubarska²⁵, E. Lohfink³⁹, C. J. Lozano Mariscal⁴¹, L. Lu³⁸, F. Lucarelli²⁸, A. Ludwig^{24,35}, W. Luszczak³⁸, Y. Lyu^{8,9}, W. Y. Ma⁵⁹, J. Madsen³⁸, K. B. M. Mahn²⁴, Y. Makino³⁸, S. Mancina³⁸, I. C. Mariş¹², R. Maruyama⁴³, K. Mase¹⁶, T. McElroy²⁵, F. McNally³⁹, J. V. Mead²², K. Meagher³⁸, A. Medina²¹, M. Meier¹⁶, S. Meighen-Berger²⁷, J. Micallef²⁴, D. Mockler¹², T. Montaruli²⁸, R. W. Moore²⁵, R. Morse³⁸, M. Moulai¹⁵, R. Naab⁵⁹, R. Nagai¹⁶, U. Naumann⁵⁸, J. Necker⁵⁹, L. V. Nguyen²⁴, H. Niederhausen²⁷, M. U. Nisa²⁴, S. C. Nowicki²⁴, D. R. Nygren⁹, A. Obertacke Pollmann⁵⁸, M. Oehler³¹, A. Olivas¹⁹, E. O'Sullivan⁵⁷, H. Pandya⁴², D. V. Pankova⁵⁶, N. Park³³, G. K. Parker⁴, E. N. Paudel⁴², L. Paul⁴⁰, C. Pérez de los Heros⁵⁷, L. Peters¹, J. Peterson³⁸, S. Philippen¹, D. Pieloth²³, S. Pieper⁵⁸, M. Pittermann³², A. Pizzuto³⁸, M. Plum⁴⁰, Y. Popovych³⁹, A. Porcelli²⁹, M. Prado Rodriguez³⁸, P. B. Price⁸, B. Pries²⁴, G. T. Przybylski⁹, C. Raab¹², A. Raissi¹⁸, M. Rameez²², K. Rawlins³, I. C. Rea²⁷, A. Rehman⁴², P. Reichherzer¹¹, R. Reimann¹, G. Renzi¹², E. Resconi²⁷, S. Reusch⁵⁹, W. Rhode²³, M. Richman⁴⁵, B. Riedel³⁸, E. J. Roberts², S. Robertson^{8,9}, G. Roellinghoff⁵², M. Rongen³⁹, C. Rott^{49,52}, T. Ruhe²³, D. Ryckbosch²⁹, D. Rysewyk Cantu²⁴, I. Safa^{14,38}, J. Saffer³², S. E. Sanchez Herrera²⁴, A. Sandrock²³, J. Sandroos³⁹, M. Santander⁵⁴, S. Sarkar⁴⁴, S. Sarkar²⁵, K. Satalecka⁵⁹, M. Scharf¹, M. Schaufel¹, H. Schieler³¹, S. Schindler²⁶, P. Schlunder²³, T. Schmidt¹⁹, A. Schneider³⁸, J. Schneider²⁶, F. G. Schröder^{31,42}, L. Schumacher²⁷, G. Schwefer¹, S. Sclafani⁴⁵, D. Seckel⁴², S. Seunarine⁴⁷, A. Sharma⁵⁷, S. Shefali³², M. Silva³⁸, B. Skrzypek¹⁴, B. Smithers⁴, R. Snihur³⁸, J. Soedingrekso²³, D. Soldin⁴², C. Spannfellner²⁷, G. M. Spiczak⁴⁷, C. Spiering^{59,61}, J. Stachurska⁵⁹, M. Stamatikos²¹, T. Stanev⁴², R. Stein⁵⁹, J. Stettner¹, A. Steuer³⁹, T. Stezelberger⁹, T. Stürwald⁵⁸, T. Stuttard²², G. W. Sullivan¹⁹, I. Taboada⁶, F. Tenholt¹¹, S. Ter-Antonyan⁷, S. Tilav⁴², F. Tischbein¹, K. Tollefson²⁴, L. Tomankova¹¹, C. Tönnis⁵³, S. Toscano¹², D. Tosi³⁸, A. Trettin⁵⁹, M. Tselengidou⁵⁷, C. F. Tung⁶, A. Turcati²⁷, R. Turcotte³¹, C. F. Turley⁵⁶, J. P. Twagirayeze²⁴, B. Ty³⁸, M. A. Unland Elorrieta⁴¹, N. Valtonen-Mattila⁵⁷, J. Vandenbroucke³⁸, N. van Eijndhoven¹³, D. Vannerom¹⁵, J. van Santen⁵⁹, S. Verpoest²⁹, M. Vraeghe²⁹, C. Walck⁵⁰, T. B. Watson⁴, C. Weaver²⁴, P. Weigel¹⁵, A. Weindl³¹, M. J. Weiss⁵⁶, J. Weldert³⁹, C. Wendt³⁸, J. Werthebach²³, M. Weyrauch³², N. Whitehorn^{24,35}, C. H. Wiebusch¹, D. R. Williams⁵⁴, M. Wolf²⁷, K. Woschnagg⁸, G. Wrede²⁶, J. Wulff¹¹, X. W. Xu⁷, Y. Xu⁵¹, J. P. Yanez²⁵, S. Yoshida¹⁶, S. Yu²⁴, T. Yuan³⁸, Z. Zhang⁵¹

¹ III. Physikalisches Institut, RWTH Aachen University, D-52056 Aachen, Germany² Department of Physics, University of Adelaide, Adelaide, 5005, Australia³ Dept. of Physics and Astronomy, University of Alaska Anchorage, 3211 Providence Dr., Anchorage, AK 99508, USA⁴ Dept. of Physics, University of Texas at Arlington, 502 Yates St., Science Hall Rm 108, Box 19059, Arlington, TX 76019, USA⁵ CTSPPS, Clark-Atlanta University, Atlanta, GA 30314, USA⁶ School of Physics and Center for Relativistic Astrophysics, Georgia Institute of Technology, Atlanta, GA 30332, USA⁷ Dept. of Physics, Southern University, Baton Rouge, LA 70813, USA⁸ Dept. of Physics, University of California, Berkeley, CA 94720, USA⁹ Lawrence Berkeley National Laboratory, Berkeley, CA 94720, USA¹⁰ Institut für Physik, Humboldt-Universität zu Berlin, D-12489 Berlin, Germany¹¹ Fakultät für Physik & Astronomie, Ruhr-Universität Bochum, D-44780 Bochum, Germany¹² Université Libre de Bruxelles, Science Faculty CP230, B-1050 Brussels, Belgium¹³ Vrije Universiteit Brussel (VUB), Dienst ELEM, B-1050 Brussels, Belgium¹⁴ Department of Physics and Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, MA 02138, USA¹⁵ Dept. of Physics, Massachusetts Institute of Technology, Cambridge, MA 02139, USA

- ¹⁶ Dept. of Physics and Institute for Global Prominent Research, Chiba University, Chiba 263-8522, Japan
¹⁷ Department of Physics, Loyola University Chicago, Chicago, IL 60660, USA
¹⁸ Dept. of Physics and Astronomy, University of Canterbury, Private Bag 4800, Christchurch, New Zealand
¹⁹ Dept. of Physics, University of Maryland, College Park, MD 20742, USA
²⁰ Dept. of Astronomy, Ohio State University, Columbus, OH 43210, USA
²¹ Dept. of Physics and Center for Cosmology and Astro-Particle Physics, Ohio State University, Columbus, OH 43210, USA
²² Niels Bohr Institute, University of Copenhagen, DK-2100 Copenhagen, Denmark
²³ Dept. of Physics, TU Dortmund University, D-44221 Dortmund, Germany
²⁴ Dept. of Physics and Astronomy, Michigan State University, East Lansing, MI 48824, USA
²⁵ Dept. of Physics, University of Alberta, Edmonton, Alberta, Canada T6G 2E1
²⁶ Erlangen Centre for Astroparticle Physics, Friedrich-Alexander-Universität Erlangen-Nürnberg, D-91058 Erlangen, Germany
²⁷ Physik-department, Technische Universität München, D-85748 Garching, Germany
²⁸ Département de physique nucléaire et corpusculaire, Université de Genève, CH-1211 Genève, Switzerland
²⁹ Dept. of Physics and Astronomy, University of Gent, B-9000 Gent, Belgium
³⁰ Dept. of Physics and Astronomy, University of California, Irvine, CA 92697, USA
³¹ Karlsruhe Institute of Technology, Institute for Astroparticle Physics, D-76021 Karlsruhe, Germany
³² Karlsruhe Institute of Technology, Institute of Experimental Particle Physics, D-76021 Karlsruhe, Germany
³³ Dept. of Physics, Engineering Physics, and Astronomy, Queen's University, Kingston, ON K7L 3N6, Canada
³⁴ Dept. of Physics and Astronomy, University of Kansas, Lawrence, KS 66045, USA
³⁵ Department of Physics and Astronomy, UCLA, Los Angeles, CA 90095, USA
³⁶ Department of Physics, Mercer University, Macon, GA 31207-0001, USA
³⁷ Dept. of Astronomy, University of Wisconsin–Madison, Madison, WI 53706, USA
³⁸ Dept. of Physics and Wisconsin IceCube Particle Astrophysics Center, University of Wisconsin–Madison, Madison, WI 53706, USA
³⁹ Institute of Physics, University of Mainz, Staudinger Weg 7, D-55099 Mainz, Germany
⁴⁰ Department of Physics, Marquette University, Milwaukee, WI, 53201, USA
⁴¹ Institut für Kernphysik, Westfälische Wilhelms-Universität Münster, D-48149 Münster, Germany
⁴² Bartol Research Institute and Dept. of Physics and Astronomy, University of Delaware, Newark, DE 19716, USA
⁴³ Dept. of Physics, Yale University, New Haven, CT 06520, USA
⁴⁴ Dept. of Physics, University of Oxford, Parks Road, Oxford OX1 3PU, UK
⁴⁵ Dept. of Physics, Drexel University, 3141 Chestnut Street, Philadelphia, PA 19104, USA
⁴⁶ Physics Department, South Dakota School of Mines and Technology, Rapid City, SD 57701, USA
⁴⁷ Dept. of Physics, University of Wisconsin, River Falls, WI 54022, USA
⁴⁸ Dept. of Physics and Astronomy, University of Rochester, Rochester, NY 14627, USA
⁴⁹ Department of Physics and Astronomy, University of Utah, Salt Lake City, UT 84112, USA
⁵⁰ Oskar Klein Centre and Dept. of Physics, Stockholm University, SE-10691 Stockholm, Sweden
⁵¹ Dept. of Physics and Astronomy, Stony Brook University, Stony Brook, NY 11794-3800, USA
⁵² Dept. of Physics, Sungkyunkwan University, Suwon 16419, Korea
⁵³ Institute of Basic Science, Sungkyunkwan University, Suwon 16419, Korea
⁵⁴ Dept. of Physics and Astronomy, University of Alabama, Tuscaloosa, AL 35487, USA
⁵⁵ Dept. of Astronomy and Astrophysics, Pennsylvania State University, University Park, PA 16802, USA
⁵⁶ Dept. of Physics, Pennsylvania State University, University Park, PA 16802, USA
⁵⁷ Dept. of Physics and Astronomy, Uppsala University, Box 516, S-75120 Uppsala, Sweden
⁵⁸ Dept. of Physics, University of Wuppertal, D-42119 Wuppertal, Germany
⁵⁹ DESY, D-15738 Zeuthen, Germany
⁶⁰ Università di Padova, I-35131 Padova, Italy
⁶¹ National Research Nuclear University, Moscow Engineering Physics Institute (MEPhI), Moscow 115409, Russia
⁶² Earthquake Research Institute, University of Tokyo, Bunkyo, Tokyo 113-0032, Japan

Acknowledgements

USA – U.S. National Science Foundation-Office of Polar Programs, U.S. National Science Foundation-Physics Division, U.S. National Science Foundation-EPSCoR, Wisconsin Alumni Research Foundation, Center for High Throughput Computing (CHTC) at the University of Wisconsin–Madison, Open Science Grid (OSG), Extreme Science and Engineering Discovery Environment (XSEDE), Frontera computing project at the Texas Advanced Computing Center, U.S. Department of Energy-National Energy Research Scientific Computing Center, Particle astrophysics research computing center at the University of Maryland, Institute for Cyber-Enabled Research at Michigan State University, and Astroparticle physics computational facility at Marquette University; Belgium – Funds for Scientific Research (FRS-FNRS and FWO), FWO Odysseus and Big Science programmes, and Belgian Federal Science Policy Office (Belspo); Germany – Bundesministerium für Bildung und Forschung (BMBF), Deutsche Forschungsgemeinschaft (DFG), Helmholtz Alliance for Astroparticle Physics (HAP), Initiative and Networking Fund of the Helmholtz Association, Deutsches Elektronen Synchrotron (DESY), and High Performance Computing cluster of the RWTH Aachen; Sweden – Swedish Research Council, Swedish Polar Research Secretariat, Swedish National Infrastructure for Computing (SNIC), and Knut and Alice Wallenberg Foundation; Australia – Australian

Research Council; Canada – Natural Sciences and Engineering Research Council of Canada, Calcul Québec, Compute Ontario, Canada Foundation for Innovation, WestGrid, and Compute Canada; Denmark – Villum Fonden and Carlsberg Foundation; New Zealand – Marsden Fund; Japan – Japan Society for Promotion of Science (JSPS) and Institute for Global Prominent Research (IGPR) of Chiba University; Korea – National Research Foundation of Korea (NRF); Switzerland – Swiss National Science Foundation (SNSF); United Kingdom – Department of Physics, University of Oxford.